

Gamma and Bessel Functions

Partial Differential Equations

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Objectives

In this lesson we will learn:

- ▶ the definition of the Gamma function,
- ▶ the properties of the Gamma function,
- ▶ how to solve Bessel's ordinary differential equation of order p ,
- ▶ notable Bessel's functions, and
- ▶ properties of Bessel's functions.

Introducing the Gamma Function

The **Gamma function** is denoted as $\Gamma(p)$ and is defined by the integral,

$$\Gamma(p+1) = \int_0^{\infty} e^{-x} x^p dx.$$

Remarks:

- ▶ The improper integral converges for all $p \geq 0$.
- ▶ For $p < 0$ the integrand becomes unbounded as $x \rightarrow 0^+$.
- ▶ The integral can be shown to converge for $p > -1$ and to diverge for $p \leq -1$.

Example

Consider the case when $p = 0$,

$$\Gamma(1) = \int_0^{\infty} e^{-x} dx = \lim_{M \rightarrow \infty} \int_0^M e^{-x} dx = \lim_{M \rightarrow \infty} (-e^{-M} + 1) = 1.$$

Connection with Factorial

Lemma

For $p > 0$,

$$\Gamma(p + 1) = p\Gamma(p).$$

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Lemma

If p is a positive integer,

$$\Gamma(p + 1) = p!.$$

Remark: since the Gamma function is defined for non-integer values of p , it generalizes the factorial.

Example: $\Gamma(1/2)$

By definition,

$$\Gamma\left(\frac{1}{2}\right) = \int_0^{\infty} e^{-x} x^{-\frac{1}{2}} dx.$$

Making the substitution $u^2 = x$ and $2u du = dx$, the integral can be re-written as

$$\Gamma\left(\frac{1}{2}\right) = \int_0^{\infty} e^{-u^2} u^{-1} (2u) du = 2 \int_0^{\infty} e^{-u^2} du = \int_{-\infty}^{\infty} e^{-u^2} du = \sqrt{\pi}.$$

Pochhammer Symbol

Lemma

For $p > 0$,

$$\frac{\Gamma(p+n)}{\Gamma(p)} = p(p+1)(p+2)\cdots(p+n-1) = (p)_n.$$

Remark: the expression $(p)_n$ is called the **Pochhammer symbol** or the **rising factorial**.

Pochhammer Symbol

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Remark: the expression $(p)_n$ is called the **Pochhammer symbol** or the **rising factorial**.

This lemma makes it possible to calculate $\Gamma(p)$ for all $p > 0$ provided $\Gamma(x)$ is known for $0 < x < 1$.

Examples

$$\Gamma\left(\frac{3}{2}\right) = \Gamma\left(\frac{1}{2} + 1\right) = \frac{1}{2}\Gamma\left(\frac{1}{2}\right) = \frac{\sqrt{\pi}}{2}$$

$$\Gamma\left(\frac{11}{2}\right) = \Gamma\left(5 + \frac{1}{2}\right) = \Gamma\left(\frac{1}{2}\right)\left(\frac{1}{2}\right)\left(\frac{3}{2}\right)\left(\frac{5}{2}\right)\left(\frac{7}{2}\right)\left(\frac{9}{2}\right) = \frac{945\sqrt{\pi}}{32}$$

Properties of the Gamma Function (1 of 2)

Lemma

The Gamma function $\Gamma(x)$ is continuous for all $x > 0$.

Properties of the Gamma Function (1 of 2)

Lemma

The Gamma function $\Gamma(x)$ is continuous for all $x > 0$.

The Gamma function is differentiable for $x > 0$.

$$\Gamma'(x) = \frac{d}{dx} \int_0^{\infty} e^{-t} t^{x-1} dt = \int_0^{\infty} \frac{d}{dx} [e^{-t} t^{x-1}] dt = \int_0^{\infty} e^{-t} (\ln t) t^{x-1} dt$$

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The Gamma function to be extended to the real numbers in the set $\{x \in \mathbb{R} \mid x \neq 0, -1, -2, \dots\}$. If n is a natural number and if $-n < x < -n+1$ then $0 < x+n < 1$ and thus $\Gamma(x+n)$ is already defined. Recall that

$$\begin{aligned}\Gamma(x+n) &= \Gamma(x)(x)_n \\ \Gamma(x) &= \frac{\Gamma(x+n)}{(x)_n}.\end{aligned}$$

This extended form of the Gamma function is also continuous on the set $\{x \in \mathbb{R} \mid x \neq 0, -1, -2, \dots\}$.

Properties of the Gamma Function (2 of 2)

Theorem

If $n = 0, -1, -2, \dots$ then $\lim_{x \rightarrow n} |\Gamma(x)| = \infty$.

Properties of the Gamma Function (2 of 2)

Theorem

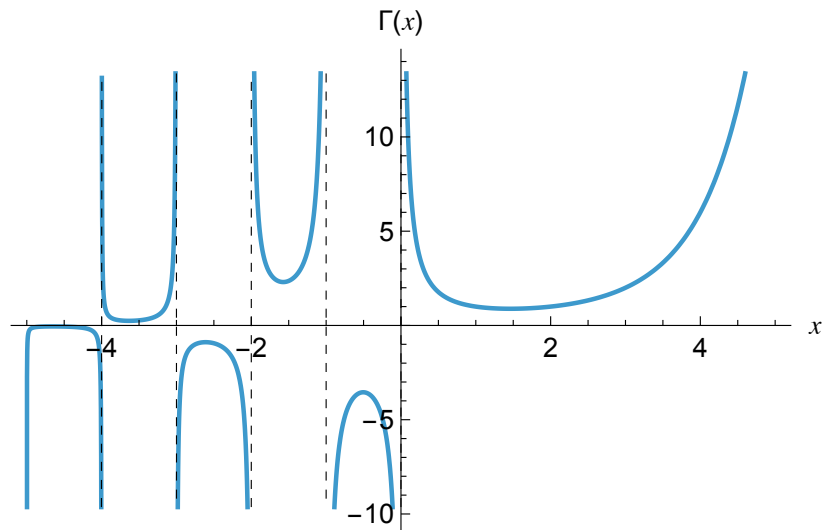
If $n = 0, -1, -2, \dots$ then $\lim_{x \rightarrow n} |\Gamma(x)| = \infty$.

- ▶ Define the reciprocal of the Gamma function at $x = 0, -1, -2, \dots$ to be 0. Thus $1/\Gamma(x)$ is continuous for all x and $1/\Gamma(x) = 0$ at $x = 0, -1, -2, \dots$
- ▶ Define $f(x) = \ln |\Gamma(x)|$ then

$$f'(x) = \frac{\Gamma'(x)}{\Gamma(x)} \iff \Gamma'(x) = \Gamma(x)f'(x).$$

- ▶ Function $f'(x)$ is called the **digamma function** or **polygamma function** and it is denoted as $\psi(x)$.
- ▶ Adopting this notation allows the derivative of the Gamma function to be written as $\Gamma'(x) = \psi(x)\Gamma(x)$.

Illustration of the Gamma Function



Bessel's Equation of Order $p \geq 0$

$$x^2 y'' + xy' + (x^2 - p^2)y = 0 \text{ for } x > 0,$$

The self-adjoint form is

$$[x y']' + \left(x - \frac{p^2}{x}\right) y = 0 \text{ for } x > 0.$$

The method of Frobenius will be used to solve this ODE.

The value $x_0 = 0$ is a regular singular point with indicial function $F(r) = r^2 - p^2$, and exponents of singularity $\pm p$.

Solving Bessel's Equation (1 of 4)

Assume $y(x) = \sum_{k=0}^{\infty} a_k x^{k+r}$, differentiate formally with respect to x ,

substitute the result into Bessel's equation and re-index and combine the series to yield

$$\begin{aligned} 0 &= x^2 \sum_{k=0}^{\infty} (k+r)(k+r-1) a_k x^{k+r-2} + x \sum_{k=0}^{\infty} (k+r) a_k x^{k+r-1} \\ &\quad + (x^2 - p^2) \sum_{k=0}^{\infty} a_k x^{k+r} \\ &= \sum_{k=0}^{\infty} [(k+r)(k+r-1) + (k+r) - p^2] a_k x^{k+r} + \sum_{k=0}^{\infty} a_k x^{k+r+2} \\ &= \sum_{k=0}^{\infty} [(k+r)^2 - p^2] a_k x^{k+r} + \sum_{k=0}^{\infty} a_k x^{k+r+2} \\ &= \sum_{k=0}^{\infty} F(r+k) a_k x^{k+r} + \sum_{k=2}^{\infty} a_{k-2} x^{k+r} \\ &= F(r) a_0 x^r + F(r+1) a_1 x^{r+1} + \sum_{k=2}^{\infty} [F(r+k) a_k + a_{k-2}] x^{k+r}. \end{aligned}$$

Solving Bessel's Equation (2 of 4)

$$0 = F(r)a_0x^r + F(r+1)a_1x^{r+1} + \sum_{k=2}^{\infty} [F(r+k)a_k + a_{k-2}]x^{k+r}$$

- ▶ Since $r = -p$ or $r = p$ then $F(r) = 0$ and for arbitrary a_0 the first term above is zero.
- ▶ The second term has coefficient $a_1F(r+1) = a_1((1+r)^2 - p^2)$ which must match the coefficient of x^{r+1} on the left-hand side of the equation. This can be achieved by choosing $a_1 = 0$ regardless of the values of p or r .
- ▶ For $k \geq 2$ the following recurrence relation holds for any r such that $F(k+r) \neq 0$,

$$a_k = -\frac{a_{k-2}}{F(k+r)}.$$

Solving Bessel's Equation (3 of 4)

$$a_k = -\frac{a_{k-2}}{F(k+r)}$$

This equation is valid as long as r is the larger of the two exponents of singularity ($r = p$). Since $a_1 = 0$ then $a_{2k-1} = 0$ for $k \in \mathbb{N}$. Since $F(k+p) \neq 0$ for $k \in \mathbb{N}$, the recurrence relation can be simplified to

$$a_k = -\frac{a_{k-2}}{k(k+2p)}.$$

Since a_0 is an arbitrary constant, then

$$\begin{aligned}a_2 &= -\frac{a_0}{2(2+2p)} = \frac{(-1)^1 a_0}{2^2(1+p)} \\a_4 &= -\frac{a_2}{4(4+2p)} = \frac{(-1)^2 a_0}{2^4 2!(1+p)(2+p)} \\&\vdots \\a_{2k} &= \frac{(-1)^k a_0}{2^{2k} k!(1+p)(2+p) \cdots (k+p)}.\end{aligned}$$

Solving Bessel's Equation (4 of 4)

One solution to Bessel's equation of order p has the form,

$$y_1(x) = a_0 \sum_{k=0}^{\infty} \frac{(-1)^k}{2^{2k} k! (1+p)(2+p) \cdots (k+p)} x^{2k+p}.$$

Traditionally the first solution to Bessel's equation of order p has the arbitrary constant a_0 chosen as

$$a_0 = \frac{1}{2^p \Gamma(p+1)}.$$

This defines the **Bessel function of the first kind of order p** :

$$J_p(x) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k! \Gamma(k+p+1)} \left(\frac{x}{2}\right)^{2k+p}.$$

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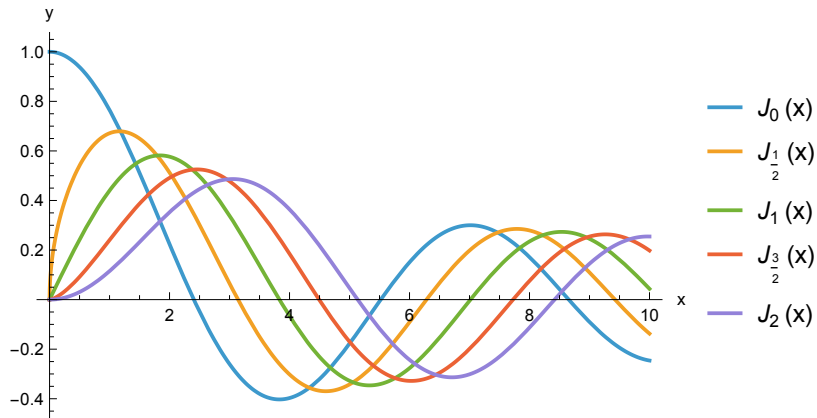
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When $p = n \in \{0, 1, \dots\}$ then $J_n(x)$ is a power series,

$$J_n(x) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k! (k+n)!} \left(\frac{x}{2}\right)^{2k+n}.$$

Bessel Functions of the First Kind



Second Linearly Independent Solution (1 of 3)

Since Bessel's equation is second order, there must be another solution that is linearly independent of $J_p(x)$.

If $2p$ is not an integer then a second solution is

$$J_{-p}(x) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k! \Gamma(k - p + 1)} \left(\frac{x}{2}\right)^{2k-p}.$$

By convention the second solution to Bessel's equation of order p , linearly independent of $J_p(x)$ is chosen to be a linear combination of $J_p(x)$ and $J_{-p}(x)$, namely

$$Y_p(x) = \frac{J_p(x) \cos(p\pi) - J_{-p}(x)}{\sin(p\pi)}.$$

The function $Y_p(x)$ is called a **Bessel function of the second kind of order p** .

Second Linearly Independent Solution (2 of 3)

If $p = n \in \mathbb{N}$ then $2p$ is an integer. Assume a second solution to Bessel's equation of order n , linearly independent from $J_n(x)$ takes the form

$$y_2(x) = aJ_n(x) \ln x + x^{-n} \left(1 + \sum_{k=1}^{\infty} c_k x^k \right)$$

where a and c_k for $k \in \mathbb{N}$ are constants.

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where a and c_k for $k \in \mathbb{N}$ are constants.

It can be shown that

$$y_2(x) = \frac{-\ln x}{2^{n-1}(n-1)!} J_n(x) + x^{-n} \left[1 + \sum_{k=1}^{n-1} \frac{(n-k-1)!}{k!(n-1)!} \left(\frac{x}{2}\right)^{2k} \right] \\ + x^{-n} \left[\sum_{k=0}^{\infty} \frac{(-1)^k (H_k + H_{n+k})}{k!(n+k)!(n-1)!} \left(\frac{x}{2}\right)^{2n+2k} \right],$$

where $H_k = \sum_{j=1}^k \frac{1}{j}$ is the k th Harmonic number.

Second Linearly Independent Solution (3 of 3)

By convention the second linearly independent solution to Bessel's equation of order $n \in \mathbb{N}$ is expressed as

$$\begin{aligned} Y_n(x) &= \frac{2}{\pi} \left((\gamma - \ln 2) J_n(x) - 2^{n-1} (n-1)! y_2(x) \right) \\ &= \frac{2}{\pi} \left(\gamma + \ln \frac{x}{2} \right) J_n(x) - \frac{1}{\pi} \sum_{k=0}^{n-1} \frac{(n-k-1)!}{k!} \left(\frac{x}{2} \right)^{2k-n} \\ &\quad - \frac{1}{\pi} \sum_{k=0}^{\infty} \frac{(-1)^k (H_k + H_{n+k})}{k!(n+k)!} \left(\frac{x}{2} \right)^{2k+n}. \end{aligned}$$

The constant γ is called the **Euler-Máscheroni constant** or merely the **Euler gamma**.

$$\gamma \equiv \lim_{n \rightarrow \infty} (H_n - \ln n) \approx 0.577216.$$

Bessel's Equation of Order $p = 0$ (1 of 2)

We have already found

$$J_0(x) = \sum_{k=0}^{\infty} \frac{(-1)^k}{(k!)^2} \left(\frac{x}{2}\right)^{2k}.$$

Assume a second solution linearly independent from $J_0(x)$ has the form

$$y_2(x) = (\ln x)J_0(x) + \sum_{k=1}^{\infty} c_k x^k,$$

where c_k for $k \in \mathbb{N}$ are constants.

It can be shown that

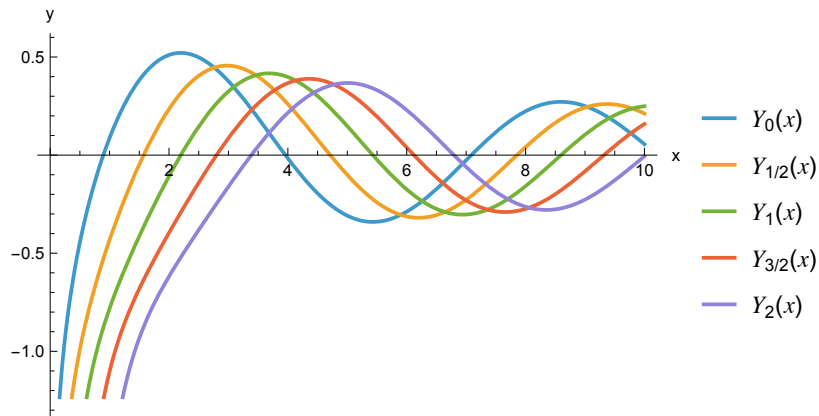
$$y_2(x) = (\ln x)J_0(x) - \sum_{k=1}^{\infty} \frac{(-1)^k H_k}{(k!)^2} \left(\frac{x}{2}\right)^{2k}.$$

Bessel's Equation of Order $p = 0$ (2 of 2)

By convention the Bessel function of the second kind of order 0 is expressed as

$$Y_0(x) = \frac{2}{\pi} \left(\gamma + \ln \frac{x}{2} \right) J_0(x) - \frac{2}{\pi} \sum_{k=1}^{\infty} \frac{(-1)^k H_k}{(k!)^2} \left(\frac{x}{2} \right)^{2k}.$$

Bessel Functions of the Second Kind



Derivative Relationships

Theorem

The Bessel functions of the first kind of order p satisfy the following derivative relationships:

$$\frac{d}{dx} [x^{-p} J_p(x)] = -x^{-p} J_{p+1}(x),$$

$$\frac{d}{dx} [x^p J_p(x)] = x^p J_{p-1}(x).$$

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Corollary

The Bessel functions of the first kind obey the following recurrence relations.

$$\begin{aligned}xJ'_p(x) - pJ_p(x) &= -xJ_{p+1}(x) \\ xJ'_p(x) + pJ_p(x) &= xJ_{p-1}(x) \\ J_{p+1}(x) + J_{p-1}(x) &= \frac{2p}{x} J_p(x)\end{aligned}$$

Zeros of Bessel Functions

Lemma

The positive zeros of $J_0(x)$ form a strictly increasing sequence which tends to positive infinity.

Theorem

Let n be a nonnegative integer and suppose $\{\lambda_{n,k}\}_{k=1}^{\infty}$ are the positive zeros of $J_n(x)$ arranged in increasing order. Between each consecutive pair of zeros of $J_n(x)$ there exists a unique zero of $J_{n+1}(x)$.

Parametric Form of Bessel's Equation

The **parametric form of Bessel's Equation** is

$$z^2 \frac{d^2 y}{dz^2} + z \frac{dy}{dz} + (\lambda^2 z^2 - p^2)y(z) = 0$$

The function $J_p(x)$ solves Bessel's equation of order p if and only if the function $J_p(\lambda z)$ solves the parametric form of Bessel's equation of order p .

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Theorem

Let $p \geq 0$ and let $\lambda_{p,k}$ be the k th zero of $J_p(x)$, then

$$\int_0^1 J_p(\lambda_{p,n} z) J_p(\lambda_{p,m} z) z \, dz = \frac{1}{2} (J_{p+1}(\lambda_{p,n}))^2 \delta_{mn},$$

for $m, n \in \mathbb{N}$, where δ_{mn} is the **Kronecker delta function**.

Bessel-Fourier Series

If $f(x)$ is sufficiently smooth on $[0, 1]$, $f(x)$ can be expressed as

$$f(x) = \sum_{n=1}^{\infty} a_n J_p(\lambda_{p,n} x),$$

where

$$a_n = \frac{2 \int_0^1 f(x) J_p(\lambda_{p,n} x) x \, dx}{(J_{p+1}(\lambda_{p,n}))^2}.$$

Bessel Functions of Half-Integer Order (1 of 3)

The **Bessel functions of half-integer order** are the functions $J_{n+\frac{1}{2}}(x)$ and $Y_{n+\frac{1}{2}}(x)$ where n is an integer.

Using the formula for $J_p(x)$ with $p = 1/2$,

$$\begin{aligned} J_{\frac{1}{2}}(x) &= \sum_{k=0}^{\infty} \frac{(-1)^k}{k! \Gamma(k+1+1/2)} \left(\frac{x}{2}\right)^{2k+\frac{1}{2}} \\ &= \sqrt{\frac{2}{x}} \sum_{k=0}^{\infty} \frac{(-1)^k x^{2k+1}}{2^{2k+1} k! \left(\frac{1}{2}\right) \left(\frac{3}{2}\right) \cdots \left(\frac{2k+1}{2}\right) \Gamma\left(\frac{1}{2}\right)} \\ &= \sqrt{\frac{2}{\pi x}} \sum_{k=0}^{\infty} \frac{(-1)^k x^{2k+1}}{(2k+1)!} = \sqrt{\frac{2}{\pi x}} \sin x. \end{aligned}$$

Bessel Functions of Half-Integer Order (2 of 3)

Using the derivative relationships between Bessel functions implies

$$\begin{aligned} J_{3/2}(x) &= \frac{1}{2x} J_{\frac{1}{2}}(x) - J'_{\frac{1}{2}}(x) \\ &= \frac{1}{2x} \sqrt{\frac{2}{\pi x}} \sin x - \sqrt{\frac{2}{\pi x}} \cos x + \frac{1}{2x} \sqrt{\frac{2}{\pi x}} \sin x \\ &= \sqrt{\frac{2}{\pi x}} \left(\frac{1}{x} \sin x - \cos x \right). \end{aligned}$$

Bessel Functions of Half-Integer Order (3 of 3)

Theorem

For $n = 0, 1, 2, \dots$

$$J_{n+\frac{1}{2}}(x) = (-1)^n x^{n+1} \sqrt{\frac{2}{\pi x}} \left(\frac{d}{x dx} \right)^n \left[\frac{\sin x}{x} \right].$$

Theorem

For $n = 0, 1, 2, \dots$

$$J_{-n-\frac{1}{2}}(x) = x^{n+1} \sqrt{\frac{2}{\pi x}} \left(\frac{d}{x dx} \right)^n \left[\frac{\cos x}{x} \right].$$

Spherical Bessel's Equation

The **spherical Bessel's equation** of order n is

$$x^2 y'' + 2x y' + (x^2 - n(n+1))y = 0.$$

The solutions are the **spherical Bessel function of the first kind**

$$j_n(x) = \sqrt{\frac{\pi}{2x}} J_{n+\frac{1}{2}}(x)$$

and the **spherical Bessel function of the second kind**

$$y_n(x) = \sqrt{\frac{\pi}{2x}} Y_{n+\frac{1}{2}}(x).$$

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Corollary

For $n = 0, 1, 2, \dots$,

$$\int_0^1 j_n(\lambda_{n+\frac{1}{2},k} x) j_n(\lambda_{n+\frac{1}{2},l} x) x^2 dx = \frac{1}{2} \left(j_{n+1}(\lambda_{n+\frac{1}{2},k}) \right)^2 \delta_{kl}$$

where δ_{kl} is the Kronecker delta function.